A SPECTROGRAPHIC STUDY OF THE INTERACTING ECLIPSING BINARY RY SCUTI: AN EPISODE IN THE RAPID MASS LOSS STAGE OR A PROTOPLANETARY NEBULA?¹

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RESUMEN

RY Scuti, la binaria de eclipse de 11 días de período, caracterizada, en su espectro, por presentar líneas de [Fe III], ha sido reinvestigada en base a material espectrográfico obtenido en Chile, en el Observatorio Europeo Austral y en el Observatorio Interamericano de Cerro Tololo, que cubre las regiones 3400–5150 Å y 5700–6700 A. Fueron también analizadas observaciones en el ultravioleta obtenidas con el satélite IUE. El espectro de RY Sct es muy complejo y peculiar, y está caracterizado por varios conjuntos de líneas de emisión y de absorción. Entre las emisiones, aparece un grupo de líneas que son características de las nebulosas planetarias. El sistema está formado por una componente B0V y una compañera que parece rodeada por una envoltura opaca que emite en He II $\lambda 4686$. En la envoltura circumbinaria podemos distinguir: (a) una región de radiación diluida animada de una velocidad de $\sim -178 \text{ km s}^{-1}$; (b) regiones donde se originan las líneas de resonancia de Si IV y C IV en el ultravioleta, con velocidades de -1200 y -600 km s⁻¹; (c) regiones que dan origen al "espectro nebular" y sugieren la existencia de una triple nebulosidad que rodea al sistema y está caracterizada por un valor de $N_{\rm e}$ probablemente mayor que 10^4 cm⁻³, un valor de $T_{\rm e} \sim 15-20,000$ K y velocidades de -18, $+9 \text{ y} + 48 \text{ km s}^{-1}$, respectivamente.

ABSTRACT

RY Scuti, the 11-day [Fe III] eclipsing binary, has been reinvestigated on the basis of spectrographic material obtained in Chile at the European Southern Observatory, La Silla, and at the Cerro Tololo Interamerican Observatory, covering the spectral regions 3400–5150 Å and 5700–6700 Å. Ultraviolet *IUE* observations were also analyzed. The spectrum of RY Sct is very complex and peculiar, being characterized by several sets of emission as well as absorption features. Among the emission features, there is a set of lines that are characteristic of planetary nebulae. The system is found to be formed by a B0V component and a companion that appears to be surrounded by an opaque envelope that emits in He II λ 4686. In the circumbinary envelope we can distinguish: (a) a region of diluted radiation that suggests a velocity of ~ -178 km s⁻¹; (b) regions of formation of the resonance lines of Si IV and C IV in the *IUE* UV, characterized by velocities of -1200 and -600 km s⁻¹; (c) regions that give rise to the "nebular spectrum" and suggest that there is a triple nebulosity surrounding the system with $N_{\rm e}$ probably larger than $10^4 \, {\rm cm}^{-3}$, $T_{\rm e} \sim 15$ -20,000 K and velocities of -18, +9 and +48 km s⁻¹, respectively.

Key Words: BINARIES: ECLIPSING — BINARIES, INDIVIDUAL: RY SCUTI — PROTOPLANETARY NEBULA

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1. INTRODUCTION

RY Scuti (HD 169515, BD -12°5045, MWC 295, SAO 161458, AFGL 5235S, IRAS 18227-1243; $\alpha_{1950} = 18^{h}22^{m}42.706; \ \delta_{1950} = -12^{\circ}43'09''24)$ is an eclipsing variable, first identified as such by Gaposchkin (1937). Its spectrum was classified as 'O9.7 Ibpe var.' by Walborn (1982). It is listed in Vol. III of the 'General Catalogue of Variable Stars' (Kholopov et al. 1987) as having visual magnitude 9.12 at maximum light, with eclipse depths $A_1 = 0.60$ mag and $A_2 = 0.44$ mag, respectively, and with the indication that the light curve varies in shape. Ciatti et al. (1980) found that the eclipses are partial, that the primary eclipse is an occultation, and that $i = 75^{\circ}$. Earlier, O'Connell (1949) had reported that the width of the primary eclipse is larger than that of the secondary eclipse, the ratio of the widths being 1.76, and that the secondary eclipse is not midway between the deeper minima but that $t_1 - t_2 - P/2 = 0.11$. For this reason, RY Sct has been listed among the systems that are candidates for displaying apsidal motion (Semeniuk 1968; Hegadiis 1969), although more recent photoelectric observations do not confirm O'Connell's findings (cf. Kumsiachvili 1985; Zakirov 1985).

The light of RY Sct is polarized and there is a considerable increase in intrinsic polarization during eclipse (Shakhovskoi 1965; Zellner & Serkowski 1972).

RY Sct is located in the galactic plane, and Banaczkowski & Dworak (1988) believe that the object is one of the most distant eclipsing binaries known in the Galaxy, at a distance of the order of 5000 pc.

Infrared observations of RY Sct have been performed in the 0.5–22 μ m interval and the list of references through 1986 may be found in the 'Catalogue of Infrared Observations' (Gezari et al. 1987). The observations indicate that the object is strongly reddened, making it the supergiant of earliest spectral type that shows the presence of dust formation, the infrared excess having the characteristics usually attributed to emission by silicate grains (Grasdalen et al. 1979). RY Sct is also an IRAS source (IRAS Catalogs and Atlases 1988), with IRAS fluxes of 44.28 Jy at 12 μ m, 22.84 Jy at 25 μ m, 22.80 Jy at 60 μ m, and 1034.86 Jy at 100 μ m, the last figure being an upper limit. RY Sct was detected by the *Einstein* satellite as an X-ray source in the 0.2–3.5 keV energy range; the upper flux limit being $\log(L_{\rm X}/{\rm erg~s^{-1}}) < 32.92$ (cf. Chlebowski et al. 1989).

Hughes & Woodsworth (1973) found that RY Sct is a radio source that seems to be of thermal nature (cf. Purton et al. 1982). Further radio observations have been reported with the VLA at 6 cm (Florkowski 1990, private communication) and with the Nancy telescope at 21 cm (Kazés 1982, private communication). However, as far as we know, these observations have not been published.

Belserene (1982, private communication) carried out a thorough analysis of the photometric observations of RY Sct available at the time, including those from the Harvard plate collection, and found that the period of the binary is "more nearly constant than had been thought". She has suggested the constant-period ephemeris

 $T_0 = JD2443342.456 + 11.124646 \,n \, days \quad (1)$

for the computation of the times of primary eclipse, which we have used throughout the paper.

From the spectroscopic point of view, RY Sct has, for a long time, attracted much interest because its spectrum is characterized by the predominance of emission features, some of which are due to [Fe III], as identified by Edlén & Swings (1939). The first attempt to determine orbital elements is due to Popper (1943), who was able to derive, from the measurements of the Si IV 4088 absorption line, a value for the mass function of 16.7 solar masses. Popper also suggested that the total mass of the system is well in excess of 100 solar masses.

A number of descriptions of the spectrum have been published in the past, concluding that there is absorption associated with some of the emission, but there was no certainty as to whether absorption lines from the secondary component were really detected. Cowley & Hutchings (1976), on the other hand, adopted the interpretation that the $H\gamma$, $H\beta$, He I 4471 and some of the other He I lines contain absorption components from both stars. They discussed several possible values of the parameters of the system and noted the very high Balmer emission decrement in the spectrum of the object.

An analysis of the optical and radio spectrum of RY Sct led King & Jameson (1979) to the conclusion that the system probably consists of a BOV primary transferring mass to a hot (\sim O5.6V) secondary. The system has been observed in the ultraviolet with the *IUE* satellite and a description of the spectrum, based on frames taken near secondary minimum, was published by Sahade et al. (1984).

2. NEW SPECTROGRAPHIC OBSERVATIONS

In view of the importance and the complexity of the RY Sct system, and in order to try to reach a better understanding of it, in 1981 a group of astronomers in the former USSR (M. B. Babaev, A. M. Cherepaschuk, M. I. Kumsiasvili, M. Yu. Skul'sky, and M. M. Zakirov) with the participation of R M. West, initiated a coordinated campaign for the observation of the object using a variety of different techniques. Later, J. Sahade, who was planning to observe RY Sct from the Cerro Tololo Interamerican Observatory, joined in, as well as E. A. Antokhina, V. V. Golovaty, V. G. Karentinov, and L. I. Snezhko. The then USSR astronomers have already published a number of papers reporting on their results (see, e.g., Bull, Abastumani Astrophysical Observatory, vol. 58, 1985). In particular, calibrated tracings of twelve ESO spectra, obtained by R. M. West in 1981 (see below), were used by M. Yu. Skul'sky for a preliminary study (Skul'sky 1985), resulting in a provisional value for the mass ratio $(M_1/M_2 = 0.3)$ which was used by Antokhina & Cherepashchuk (1988), in connection with the photometric data from Ciatti et al. (1980), Kumsiasvili (1985), and Zakirov (1985), to derive photometric elements, in particular $i = 84^{\circ}.6$. Golovaty & Skul'sky (1990) also used the above mentioned ESA spectra, supplemented by 27 Å $\rm mm^{-1}$ spectra from the sixmeter BTA telescope in the former USSR, to study the physical characteristics of the nebular envelope around RY Sct.

The present investigation is based on a new, complete set of ground-based and *IUE* spectra, the former with higher resolution and better phase coverage than had been attained earlier. They were obtained by two of us (R. M. West and J. Sahade) at the European Southern Observatory, on Cerro La Silla, in Chile, with the coudé spectrograph attached to the 152 cm reflector, yielding dispersions of about 12 Å mm^{-1} (22 spectra) and about 20 Å mm^{-1} (2 spectra), and by one of us (J. Sahade) at the Cerro Tololo Interamerican Observatory with a coudé spectrograph attached to the 152 cm reflector and giving dispersions of $\sim 9 \text{ Å mm}^{-1}$ in the blue and ~ 18 Å mm⁻¹ in the red (9 spectra) and \sim 18 Å $\rm mm^{-1}$ in the blue and \sim 36 Å $\rm mm^{-1}$ in the red (3 spectra). The spectra cover the regions 3400–5150 Å and 5700–6700 Å.

All spectra were scanned on the PDS measuring machine at the ESO headquarters and reduced by means of the HIAP image processing system installed there. A $10 \times 10 \ \mu m$ slit was used and the photometric calibration was made by means of spot sensitometer marks (Tololo) or sensitometric spectra from the echelec facility (ESO). The wavelength scale was established by means of the Fe arc comparison spectra, the mean r.m.s. being ~ 0.2 Å. The observing log is given in Table 1; the different columns containing: (a) the identification of the spectra on the ESO and CTIO archives; (b) the Heliocentric Julian Dates of mid-exposure (HJD); (c) the corresponding orbital phase derived from equation (1); (d) the exposure time in minutes; (e) the spectral range used for the present investigation; f) the S/N ratio as measured in two regions (4030–4060 Å and 4650–4675 Å) where only few lines are present; and (g) the running number when the spectra are arranged according to phase.

Ten IUE spectra in the low dispersion, large aperture mode, were obtained by one of us (R. West) at the ESA VILSPA ground station. They were supplemented by 8 spectra from the IUE archives that were originally taken by S. Sobieski, J. Sahade and A. Michalitsianos, also in the low dispersion, large aperture mode. In total, we had at our disposal 8 images taken with the SWP camera (1165–2125 Å) and 10 images from the LWR camera (1845–3230 Å).

3. THE ULTRAVIOLET SPECTRUM

According to Sahade et al. (1984), the ultraviolet (*IUE*) spectrum of RY Sct corresponds to an object with $T_{\rm e} = 30,000$ K (spectral type B0) and $\log g = 4.0$. The line identifications by these authors are, in general, confirmed, the most important addition referring to the presence of NV in emission; on some of the spectra, these lines appear to have P Cyg-type profiles. This is certainly the case for the resonance lines of Si IV and C IV, where the P Cyg absorptions indicate velocities of -1200 and -600 km s⁻¹, respectively.

A careful examination of the SWP spectra did not reveal any changes with phase. As for the LWR spectra, they are noisy and do not allow us to ascertain whether the Mg II feature at 2800 Å is actually absent at primary eclipse, as they would seem to suggest.

Table 2 contains the updated list of line identifications in the *IUE* wavelength range.

4. THE PHOTOGRAPHIC SPECTRUM

The spectrum of RY Sct is rich in lines, very complex and unusually peculiar. It is characterized by the presence of emission as well as absorption features that blend together, making the derivation of orbital radial velocities from them extremely difficult. Table 3 lists the elements with a couple of relevant comments.

Identifier	HJD $(2,440,000+)$	Phase	Exposure ^a	Range $(Å)$	S/N	No.
ESO Spectr	ra					
G11990	4775.7554	0.840	281	3470-5095	21	25
G11991	4776.7002	0.925	300	3470 - 5095	9	29
G11992	4777.6460	0.010	300	3470 - 5095	19	1
G11993	4777.8155	0.025	120	3850 - 4915	8	2
G11998	4778.5818	0.094	111	3470 - 5095	21	3
G11999	4778.6926	0.104	192	3470 - 5095	25	4
G12000	4778.8231	0.116	164	3470 - 5095	19	5
G12004	4779.6252	0.188	206	3470 - 5095	23	8
G12005	4779.7585	0.200	152	3470 - 5095	24	9
G12006	4779.8599	0.209	120	3470 - 5095	14	11
G12010	4780.5950	0.275	169	3470 - 5095	24	14
G12011	4780.7252	0.287	188	3470 - 5095	25	15
G12579	5119.6822	0.756	150	3470 - 5095	16	21
G12580	5119.7853	0.765	133	3470 - 4755	23	22
G12581	5119.8891	0.774	150	3470 - 5095	18	23
G12587	5120.6916	0.846	141	3670-5060	16	26
G12588	5120.8190	0.858	210	3470 - 5095	24	27
G12594	5123.7261	0.119	180	3470 - 5095	21	6
G12595	5123.8594	0.131	186	3470 - 5095	25	7
G12600	5124.7011	0.207	236	3470-5010	24	10
G12601	5124.8498	0.220	136	3470-5010	17	12
F07994	5182.6034	0.412	225	5750-6685		
G12664	5183.5451	0.496	194	3470-5050	17	18
F08001	5184.5311	0.585	172	5750-6685		
CTIO Spec	$\mathrm{tra}^{\mathrm{b}}$					
D01237	4802.5846	0.252	214		7	13
D01241	4803.5547	0.339	180	3470 - 5095	16	16
D01249	4804.7957	0.450	249		2	17
D01252	4805.7283	0.534	360	3615 - 4900	11	19
D01257	4807.5997	0.703	391	3850 - 4885	10	20
D01260	4808.6024	0.793	335	3470-4940	11	24
D01263	4809.5344	0.876	235	3470-4940	13	28

 TABLE 1

 OBSERVATION LOG: SPECTROGRAPHIC OBSERVATIONS

^aIn minutes. ^bThe red plates (5750–6685 Å) were obtained simultaneously; they are all very good.

Wavelength (Å)	Character ^a	Identification
1240.6	Е	N V(1)
1258	А	S II(1) 53.79, 59.53, 50.50;
		Fe II(9) 60.54 ; SiII(4) 64.74 , 60.42
1299.2	А	Si III(1) 98.66, 94.67
1333	А	C II(1) 34.5323, 35.6627, 35.1077
1390.3	А	Si IV(1) 93.755
1393	\mathbf{E}	Si IV(1) 93.755
1398	А	Si IV(1) 02.770
1403	\mathbf{E}	Si IV(1) 02.770
1422	А	C III() 20.033, 20.438, 21.640, 21.755, 22.408
1454	А	C II(5) 55.194
1494	А	C II(3) 98.697
1532	А	P II(1) 32.51, 35.90, 36.39, 42.29, 43.09, 43.61
1545	А	C IV(1) 48.185, 50.774
1553	\mathbf{E}	C IV(1) 48.185, 50.774
1588	\mathbf{E}	[Ne V] 74.8
1598 - 1633	А	Si III(1) 98.002, 3.064, 10.094; Fe (8)
1654	А	Fe II $(40,41)$ 59.49, 58.78
1669	А	Al II(2) 70.7867
1694	А	Fe II (38) 02.043
1710 - 32	А	Fe II (38) 12.977, 20.616; Al II (6)
		19.4400, 21.2436, 21.2714, 26.39
1851	А	Al III(1) 54.72, 62.78; Al II(4) 55.93, 62.31
2338	А	Fe II (3) 32.798, 43.495, 38.005, 44.278
2362	А	Fe II (3) 59.11, 64.82
2391	А	${\rm Fe}{\rm II}(2)95.63,82.03,88.63,99.241,99.237$
2410	А	Fe II(2) 04.88, 10.52, 13.31, 06.66, 11.06
2580^{b}	А	Fe II(1) 85.88 , 98.37 , 99.40 , 07.09 and also 11.87 ,
		13.82,17.62,21.67,25.66,28.29,31.05,31.32
2820	А	Mg I(1) 52.127 (1000)

TABLE 2THE IUE SPECTRUM OF RY SCUTI

^aA: absorption; E: emission ^bPossibly partly superimposed on *reseau* mark.

TABLE 3 ELEMENTS PRESENT IN THE SPECTRUM OF RY SCUTI

Element/Ion
Ha
$\mathrm{He}\mathrm{I},\mathrm{He}\mathrm{II}^\mathrm{b}$
C II, C III
N II, [N II], N III
[O I]?, O II, [O II], [O III]
Ne II, [Ne III]?
${ m MgII}$
Al III
Si II?, Si III, Si IV
S III, [S III]?
Ca I, Ca II
Mn II
[Fe III]
[Ni III]

^aAt least up to H24, probably up to the Balmer limit. ^bHe II only as a broad emission at λ 4686 Å.

4.1. The Emission Spectrum

Four major systems can be distinguished in the emission spectrum:

(a) Very strong, relatively narrow features that show in the nebular lines: of [N II] 6766, 6549 (Fig. 1) and 6584; [O III] 4383 (very weak), 4959 and 5007; [O II] 3726, 3729; in the transitions of [Fe III], [O I], [S III] and [Ni III], and in the permitted transitions of H (Fig. 1 also shows $H\alpha$), He and Al III. In what follows we will refer to these lines as the E1 emission system.

(b) A set of broad hydrogen emissions (E2) (cf. Fig. 2, redward of the emission referred to in (a)), apparently part of broad P Cyg profile.

(c) A broad, rather faint emission at 4560 Å (E3) (Fig. 3), so far unidentified, present in symbiotic and other emission-line objects (cf. Meinel, Aveni & Stockton 1968)

(d) A broad emission of He II 4686 (E4) (Fig. 4).

4.1.1. The E1 Emission System (cf. Table 4)

The E1 emission system comprises three distinct velocity components in H α (Fig. 5) and in H β (Fig. 6). The components in H α , found by means of iterative Gaussian profile-fitting, have velocities +46.4, +6.8 and -16.0 km s⁻¹, respectively. In H β , the corresponding velocities are +48.9, +11.2 and



Fig. 1. The region of [N II] λ 6549 and H α in the spectrum of RY Sct. The ordinate scale is arbitrary, and the figures on the abscissae are Å.



Fig. 2. The region of $H\beta$ in the spectrum of RY Sct. The ordinate scale is arbitrary, and the figures on the abscissae are Å.



Fig. 3. The feature at λ 4560 in the espectrum of RY Sct. The ordinate scale is arbitrary, and the figures on the abscissae are Å.



Fig. 4. The region of the broad emission of He II λ 4686 Å in the spectrum of RY Sct, throughout the orbital cycle. Strip number 13 corresponds to phase 0.252. Phases increase from bottom upwards.

-18.3 km s⁻¹. In neither of these lines could a good fit be obtained with two Gaussians only, and the near-perfect three-component fits make the presence of further components unlikely. In H γ (Fig. 7), the three components are still there, but it was possible to separate the two most violet ones only if their velocities were assumed equal to the mean of those measured at H α and H β . If we designate the three components as a, b and c, in order of decreasing velocity, their mean equivalent widths, in Å and their mean velocity $\langle V \rangle$ in km s⁻¹, are as given in Table 6.

Most of the remaining strong, relatively narrow emission lines appear double. The mean values of the radial velocities from the two components are $+48.4\pm1.0$ km s⁻¹ (32 lines) and $+3.9\pm1.4$ km s⁻¹ (32 lines), respectively, indicating that the shortward component is actually a blend of the two violet components (b and c) that are definitely present in the first members of the Balmer series.

The E1 emissions do not appear to vary in velocity with the phase of the orbital motion. The equivalent widths of the nebular lines and the radial velocities of the components, measured on the mean spectrum, are also listed in Table 4. The total equivalent widths agree reasonably well with those of Golovaty & Skul'sky (1990), although these authors do not separate the individual components. The [N II] lines that are in the neighborhood of H α are shown in Figs. 8 and 9. The spectral range of our material

TABLE 5 THREE-GAUSSIAN FITS TO H β , H α AND [N II] 6584

	Peak Wavelength (Å)	Radial Velocity $(\rm km~s^{-1})$
${ m H}eta$	4861.05	-17.5
	4861.50	+10.2
	4862.11	+47.8
$H\alpha$	6562.29	-23.0
	6562.89	+3.5
	6563.84	+46.8
[N II]	6583.30	-13.5
	6583.80	+ 9.0
	6584.63	+47.0

THREE-GAUSSIAN FITS TO BALMER LINES

	Equivε	$\langle V \rangle$		
Component	$H\alpha$	${ m H}eta$	$ m H\gamma$	$(\rm km~s^{-1})$
a	3.600	0.862	0.263	+47.7
b	1.782	0.477	0.131	+ 9.0
с	1.134	0.342	0.100	-17.6

covers the regions of the forbidden lines [O II] at 3729 Å and 3727 Å and [N II] at 5755 Å and 6584 Å the ratios of which can provide the values of $T_{\rm e}$ and $N_{\rm e}$ that characterize the nebulosity responsible for them (cf. Gurzadyan 1969, Osterbrock 1989, Eissner et al. 1969, Seaton & Osterbrock 1957, Nussbaumer 1971). However, in our material the [O II] pair is faint and the reality of the feature at 3729 Å is doubtful. As a result, from the present material one can only say that $N_{\rm e}$ is most probably greater than $10^4 {\rm ~cm^{-3}}$ and that $T_{\rm e} \sim 15,000-20,000 {\rm ~K}.$

On the other hand, the Balmer decrement is steeper in RY Sct than in planetary nebulae and, as a consequence, the plot of $\log[F(H\gamma)/F(H\beta)]$ versus $\log[F(H\alpha)/F(H\beta)]$ falls outside the theoretical predictions as well as outside the empirical distribution of values for planetary nebulae that are shown in Barker (1974). The departures are in the sense of a steeper Balmer decrement, and consequently a smaller density (cf. Miyamoto 1952), in RY Sct than in planetary nebulae. The ratios between the intensities of the strongest components in [O III] 5007, [O III] 4959 and H β , 0.3:0.009:1, are in the range of values found in diffuse nebulae but not in planetary

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STRONG, NARROW EMISSION FEATURES IN THE SPECTRUM OF RY SCUTI ^a						
	Radial Velocity (km s ^{-1})					
Ion	$\lambda_0~({ m \AA})$	λ_{obs} (Å)	EW (Å)	single-peaked	doub	le-peaked
H 14	3721.940	21.90			+3.2	
		22.30				+31.4
[O II]	3726.16	26.60	0.022	+35.4		
[O II]	3728.91	28.80	0.019	-8.8		
H 12	3750.154	50.15			-0.3	
		50.84				+54.9
H 11	3770.632	70.65			+1.4	
		71.30				+53.1
H 10	3797.900	98.05			+11.8	
		98.58				+51.3
${\rm HeI}$	3819.697	20.35		+42.1		
H 9	3835.386	35.35			+2.8	
		36.00				+48.2
${ m HeI}$	3888.646	89.30		+50.5		
$H\epsilon$	3970.074	70.15			+5.7	
		70.70				+47.3
[Fe III]	4008.3	08.35		+3.7		
${ m HeI}$	4026.218	26.30			+6.1	
		26.90				+50.8
$[{\rm FeIII}]$	4046.4	45.75		-48.2		
$[{\rm FeIII}]$	4079.7	80.60		+66.2		
[Fe III]	4096.6	96.01		-43.2		
${ m H}\delta$	4101.737	02.40		+48.5		
HeI	4120.857	20.50			-26.0	
		21.40				+39.5
[Fe III]	4144.3	43.85		-32.6		
			0.100			
$ m H\gamma$	4340.468	40.55	0.131		+5.7	
		41.20	0.263			+50.6
HeI	4387.928	87.90			+1.9	
		88.75				+56.2
HeI	4471.508	71.60	0.219		+6.2	
		72.25	0.375			+49.8
[Fe III]	4607.0	07.10			+6.5	
		07.85				+55.4
[Fe III]	4658.1	58.20	0.467		+6.4	
_		58.90	0.584			+51.5
[Fe III]	4701.5	01.60	0.259		+6.4	
		02.40	0.338			+57.4
[Fe III]	4733.9	34.00			+6.4	
		34.65				+47.5

TABLE 4

	Radial Velocity (km s^{-1})					(s^{-1})
Ion	λ_0 (Å)	$\lambda_{ m obs}$ (Å)	EW (Å)	single-peaked	dou	ble-peaked
[Fe III]	4754.7	54.90			+12.6	
		55.50				+50.5
[Fe III]	4769.4	69.55			+9.4	
		70.20				+50.3
[Fe III]	4777.7	77.65			-3.1	
		78.45				+47.1
[Fe III]	4813.9	14.35		+28.0		
			0.342			
${ m H}eta$	4861.332	61.56	0.477		+14.1	
		62.10	0.862			+47.4
[Fe III]	4881.0	81.00			0	
		81.75				+46.1
HeI	4921.929	22.75		+50.0		
[Fe III]	4930.5	31.45		+57.8		
[O III]N	4958.91	58.80	0.054		-6.7	
10 1		59.75	0.079			+50.8
[O III]N	5006.84	06.85	0.244		+0.6	
[T]	5011.0	07.65	0.272		2.0	+48.5
[Fe III]	5011.3	11.25			-3.0	
[I] 7 7 7		11.95			F 1	+38.9
[Fe III]	5047.736	47.65			-5.1	
TT T		48.39				+38.7
Hel	5015.675	15.75			+4.5	
		16.50	0.017		150	+49.3
	5754.8	54.90	0.217		+5.2	. 41 7
II. I	EOTE CCC	55.00 75.90	0.302			+41.7
men	3873.000	75.60 76.60			+0.8	177
$[\Omega I]$	6300 23	70.00		133		+41.1
	6312.1	12.26		± 0.0	± 7.6	
	0312.1	12.20			± 1.0	±46.6
[N; III]	6401 5	02.00		± 23.4		40.0
	0401.0	02.00	0.747	120.4		
[N 11]	6548 1	48.39	0 7149		+13.1	
[1,11]	0010.1	49.20	1 046		1 10.1	+50.4
$H\alpha$	6562.817	10.20	1.134			100.1
1100	0002.011	63.10	1.782		+12.9	
		63.90	3.660		1 12:0	+49.5
[N II]N	6583.6	00.00	2.136			1 2010
[-·-]+·		83.80	0.974		+9.1	
		84.60	3.512			+45.6
He I	6678.149	78.40			+11.3	
		79.40				+56.2

TABLE 4 (CONTINUED)

^aThe peak wavelengths and corresponding radial velocities of the triple Gaussian–fitted profiles of H β , H α , and N II 6584 are given in Table 5.



Fig. 5. The H α line profile in the spectrum of RY Sct as the sum of 3 components. The ordinate scale is arbitrary, and the figures on the abscissae are Å.



Fig. 6. The H β line profile in the spectrum of RY Sct as the sum of 3 components. The ordinate scale is arbitrary, and the figures on the abscissae are Å.

nebulae (cf. Gurzadyan 1969). Moreover, the ratios measured in RY Sct,

$$I([O III]5007 + [O III]4959)/I(H\beta) = 0.4,$$

and

$$I([N \text{ II}]5755 + [N \text{ II}]6584)/I(\text{H}\beta) = 5,$$

are also outside the values that are typical of planetary nebulae: +1.8, -20.2 and -1, respectively (Gurzadyan 1969).

4.1.2. The E2 Emission System (cf. Table 4)

The broad E2 emission system that can be seen in Figs. 5 and 6 does not appear to display significant variations in velocity with phase.

4.1.3. The E3 Emission System

As can be seen in Fig. 3, the velocity of the broad emission at about 4560 Å appears to follow the same trend as Si IV 4088 (§ 4.2.2) and H9 (§ 4.2.1), but with a smaller amplitude, of the order of -67 km s^{-1} . It should be noted that the feature is bordered and certainly affected by the absorptions of Si III 4552



Fig. 7. The H γ line profile in the spectrum of RY Sct where the two most violet components (1 and 2) that are detected at H α and H β , cannot be distinguished separately in this case. The ordinate scale is arbitrary, and the figures on the abscissae are Å.



Fig. 8. Components in the [N II] λ 6548 line of the spectrum of RY Sct. The ordinate scale is arbitrary, and the figures on the abscissae are Å.



Fig. 9. Components in the [N II] $\lambda 6584$ line of the spectrum of RY Sct. The ordinate scale is arbitrary, and the figures on the abscissae are Å.

and of Si III 4567. The apparent shift of λ 4560 Å might, therefore, be thought to be the result of the shifts of these absorptions. Indeed, Skul'sky (1985) based his original solution of the orbital elements on the apparent shift of the Si III line, which may arise in the two components of the system. The central part of the 4650 Å emission profile, which is unaffected by the absorption lines, clearly moves, but because of the low S/N ratio and the shallowness of the feature, both the Gaussian fitting and the cross-correlation techniques gave rather divergent values.

4.1.4. The E4 Emission System

Cross-correlation of the individual spectra against the mean spectrum in the spectral interval 4670–4700 Å as well as the alignment of the spectra in order of phase (Fig. 4), strongly suggest that the gaseous mass that produces the He II 4686 Å emission line moves with the secondary star of the system, if we choose to call "secondary" the star that is in front at principal eclipse. The amplitude is of the order of 150 km s⁻¹, that is, significantly smaller than that of Si IV 4088, in agreement with the finding of Cowling & Hutchings (1976).

4.2. The Absorption Spectrum

In the absorption spectrum, it is possible to distinguish:

(a) A set of H and He I lines which has a total width of 500–600 km s⁻¹ (cf. Fig. 10). The He I lines, which are present throughout the entire spectrum, and the H lines are always distorted by super-imposed emission and narrow absorptions, as may be visualized in Fig. 10 in the case of H δ ; in what follows these will be referred to as the A1 absorption system.

(b) A set (A2) of much narrower lines, 250-300 km s⁻¹ in total width, as measured in the case of Si IV 4088 (Fig. 11), displayed by the rest of the elements present in the spectrum (cf. Table 4).

(c) Strong, narrow absorption features (A3) that are located to the violet of a narrow, strong emission at about λ 3890 Å (Fig. 12).

(d) Strong interstellar features (A4), some of which are also very broad, their wavelengths being listed in Table 7. The Si III profiles at 4552.7 Å must be somewhat distorted by a broad emission that appears to move together with the absorption but with smaller amplitude (cf. Fig. 13).



Fig. 10. The region of $H\gamma$ in the spectrum of RY Sct. The ordinate scale is arbitrary, and the figures on the abscissae are Å.



Fig. 11. The region of the absorption line of Si IV 4088 Å in the spectrum of RY Sct throughout the orbital cycle. Strip numbered 14 corresponds to phase 0.275. Phases increase from bottom upwards.



Fig. 12. The region of H I 3888 in the spectrum of RY Sct. The ordinate scale is arbitrary, and the figures on the abscissae are Å.

	$\lambda_{ ext{measured}}$		$\mathrm{FWHM^{b}}$	$\mathop{\mathrm{EW}}_{\circ}$	Radial Ve	locity
Identification ^a	(A)	Comments	(A)	(\mathbf{A})	$(\rm km \ s^{-}$	·1)
Ca II-K	3933.60		0.601	0.606	-4.9	+71.7
Ca II-H	3968.40		0.614		-5.3	+70.3
CaI	4226.80			0.038	+5.1	
CH II	4232.50			0.028		
CH I	4300.35			0.038		
H_2	4428.20	very broad	22.26°			
?	4762.60					
$Na I-D_2$	5889.80		1.30	0.912	-7.8	
$\operatorname{Na} I-D_1$	5896.00		1.11	0.966	+3.9	
?	6203.60					
?	6270.4					
?	6284.2					
?	6613.6	very broad	8.3°	2.09^{c}		
Mean value					-1.8 ± 1.34	+71.0

TABLE 7 INTERSTELLAR LINES IN THE SPECTRUM OF BY SCUTI

^bMeasured on a mean spectrum. ^acf. Scheffler(1982)

^cValues derived through a Gaussian fit to the line.



Fig. 13. The orbits of the components of RY Sct. The triangles correspond to the measurements of the HeII 4686 Å absorption line and the crosses correspond to the measurements of the emission feature of Si IV 4088 Å.

4.2.1. The A1 Absorption System

Because of the complex line profiles and low S/N ratio, it is impossible to directly measure meaningful orbital radial velocities from any of the broad stellar absorptions in the spectrum of RY Sct. Therefore, we decided to attempt the determination of a velocity curve by the application of the cross-correlation method to the spectral region where the distortion of the lines appear to be the smallest, i.e., to the 3730–3845 Å interval. This spectral region contains four Balmer lines (H9–H12) and one helium line (He I 3820) that are strong, as well as some much weaker lines of different ions. After truncation of this spectral region, each of the spectra was cross-correlated against the same region in the averages of the spectra taken at phases 0.104 and 0.116 (spectra numbers 4 and 5 in Table 1), and also against the average of the spectra taken at phases 0.200 and 0.207 (spectra numbers 9 and 10 in Table 1). However, this approach did not lead to any reliable determination of the shifts, probably because the lines at the violet end of the region are relatively weaker and because in He I 3820 the E1 emission is comparatively strong. H9 would remain as the only line suitable for radial velocity measurement, although the profile does not appear to be simple.

4.2.2. The A2 Absorption System

It might be thought that the A2 lines are actually the same as the A1 lines and that the difference in total width is only the result of a blending effect, perhaps with broad emissions. However, this cannot be so—if it were, then these lines, e.g., Si IV 4088, would have a total width of over 700 km s⁻¹. This would, in any case, make them different from the A1 features.

The reality of the presence of two sets of absorption lines that behave differently, in what concerns both the velocity and the width, is clearly illustrated

in Fig. 11; during the second half of the orbital cycle, the N III 4103 absorption displays a longward shift similar to that of Si IV 4088, and it is definitely larger than that of H δ .

Among the A2 lines, we select Si IV 4088 for measurement, i.e., the same line that Popper (1943) used to determine the velocity curve of the primary component of the system. This was done by fitting the individual spectra with suitable Gaussian profiles; the radial velocities thus derived are also presented in Table 8.

4.2.3. The A3 Absorption System (Fig. 12)

The spectral profile near $\lambda 3888-90$ displays the strong E1 emission as well as two strong, deep absorptions. The most violet-displaced absorption should obviously be identified with He I 3889; and this identification yields a velocity of -178 km s⁻¹. As for the less violet-displaced feature that corresponds to a wavelength of 3888 Å it must be due to H8 with some He I blending.

4.2.4. The A4 Absorption System and the Distance to RY Sct

The interstellar lines were identified and measured in the less noisy, mean spectrum of RY Sct, and are listed in Table 7 together with their equivalent widths and radial velocities. The average velocity from five interstellar lines is -1.8 ± 1.34 km s⁻¹. The H and K lines of Ca II display an additional, relatively weak, component at +71 km s⁻¹.

The measured equivalent width of the strongest interstellar Ca II-K component, EW = 0.61 Å when inserted into Evans (1941) relation, gives the distance d = 1.8 kpc. Moreover, the strengths of the D lines of Na I yield (cf. Scheffler 1982) a distance of 2.2 kpc. We, therefore, conclude that RY Sct is at a distance of ~ 2 kpc, which implies an absolute magnitude $M_V = -2.4$ for the system.

The derived value for the distance of RY Sct is in reasonable agreement with the value of 1.6 kpc derived by Milano et al. (1981) on the basis of the star's color excess. Moreover, Zakirov (1985), from photoelectric photometry, concludes that RY Sct is a member of the Ser OB I association, the distance of which is given as d = 1.5 kpc by Lindoff (1968).

All the above mentioned conclusions regarding the distance of RY Sct definitely rule out Banaczkowski & Dworak's figure of 5 kpc for the actual distance of RY Sct.

5. THE VELOCITY CURVES

In summary, there are four sets of lines in the spectrum of RY Sct that behave as if they indicate orbital motion:

- (i) E4: the broad emission of He II 4686, with a total amplitude of $\sim 150 \text{ km s}^{-1}$;
- (ii) A2: the Si IV 4088 and similarly behaved lines, with an amplitude of \sim 530 km s⁻¹;
- (iii) A1: the H and He I broad absorptions, with an amplitude of $\sim 170 \text{ km s}^{-1}$;
- (iv) E3: the emission at 4560 Å (cf. 4.1.3).

The results from He II 4686 suggest, as the alignment of the spectra had already done (Fig. 4), and as mentioned in § 4.1.4, that this emission arises in a thick envelope that surrounds the star, which is in front at primary eclipse and partakes of its orbital motion. This feature cannot be interpreted in any other way because of the profile displayed as well as because of the velocities involved. Therefore, the He II 4686 line provides information regarding the orbit of the secondary component of the system. Note, however, that Fig. 4 suggests that the violet side of the line is distorted by a neighboring absorption line and, therefore, the actual value of the velocities could be somewhat larger than those in Table 8.

The existence of some kind of disk around the secondary component of the system indicates that in all probability it is the star that we call the primary component that is losing mass in the system of RY Sct.

We now describe the absorption lines that undergo velocity variations in phase with the orbital motion of the primary component, i.e., the A2 and A1 absorption systems.

The A2 set of lines appears to show no distortion of any kind at any phase, while the A1 set of lines, represented by H9, are broader and yield very large negative velocities in the phase interval 0.094–0.131. These large approach velocities may be interpreted as arising in a mass of gas streaming from the primary component towards the companion, as is the case in many interacting binaries (cf. β Lyrae, for instance). It would be natural to think that Si IV 4088 and the similarly behaved lines (A2) arise in the photosphere of the primary star and reflect the motion on its orbit around the center of gravity of the system. This is supported by the fact that the relevant lines have the relative intensities that correspond to a B0 star.

The H9 line gives a velocity curve with a much smaller amplitude than Si IV, but this is a common, though not yet completely understood, feature in

TABLE 8

RADIAL VELOCITIES USED FOR THE ORBITAL SOLUTION OF RY SCT

Phase	Radial	Velocity ^a
(P)	Si IV 4088	He II 4686
0.010		-77.6
0.025	+45.9	-123.8
0.094	+43.3	-182.8
0.104	+42.1	-166.8
0.116	+31.0	-202.4
0.119	+57.1	-169.7
0.131	+34.0	-202.4
0.188	+89.0	-214.1
0.200	+92.0	-195.4
0.207	+108.8	-185.0
0.209	+111.0	-216.0
0.220	+139.2	-212.4
0.252	+165.0	-231.0
0.275	+112.8	-199.1
0.287	+78.3	-240.7
0.339	+112.2	-180.3
0.450	+8.9	-85.3
0.496	+40.4	-20.2
0.534	-19.0	+3.6
0.702		+244.8
0.756	-46.4	+338.2
0.765	+10.7	
0.774	-7.1	+286.1
0.793	-58.6	+304.4
0.840	-17.5	+171.5
0.846	-35.6	+210.6
0.858	-70.0	+145.2
0.876	+43.4	+127.3
0.925	+26.8	+77.6

^aUnits: km s⁻¹.

systems whose velocity curves are derived from spectral lines arising in the matter of the gaseous envelope. The above considerations suggest that the orbital elements of the RY Sct system should be derived from a simultaneous solution of the velocities of Si IV 4088 and of He II 4686 (Table 8). This was carried out with the use of the Lehmann-Filhés method and the resulting orbital parameters are listed in Table 9 and illustrated in Fig. 13. With $i = 84^{\circ}.6$ (cf. § 2), we obtain masses of about $10 M_{\odot}$ for the primary B0 star and $36 M_{\odot}$ for the companion. The

TABLE 9

ORBITAL ELEMENTS OF THE RY SCUTI SYSTEM

p	11.1246		days
T	$\rm JD45107744$		
e	0.16	± 0.04	
ω	344.8	± 15.4	degrees
$V_{\rm o}$	17.4	± 5.0	${\rm km}~{\rm s}^{-1}$
K_1	272.8	± 13.5	${\rm km~s^{-1}}$
K_2	70.8	± 8.9	${\rm km~s^{-1}}$
$a_1 \sin i$	$4.1 imes 10^7$	$\pm 2.06 \times 10^6$	km
$a_2 \sin i$	$1.1 imes 10^7$	$\pm 1.35 \times 10^6$	km
$M_1 \sin^3 i$	9.3	± 1.5	M_{\odot}
$M_2 \sin^3 i$	35.8	± 3.8	M_{\odot}

derived values of the system elements are given in Table 8. Fig. 13 suggests that phase 0 in Belserene's photometric elements does not correspond to actual conjunction.

In conclusion—and summarizing the paper—we can say that RY Sct is an interacting binary formed by a B0 primary of $\simeq 10 M_{\odot}$ and a secondary of $\sim 36 M_{\odot}$, which is surrounded by a thick envelope or disk that emits in He II 4686. The rotational velocity of the B0 star is of the order of 120–150 km s⁻¹. The system is caught when the star that is losing mass is already the less massive component, that is, after the mass-ratio reversal has taken place.

The system is embedded in an extended envelope where we can distinguish the following layers:

(a) A region dominated by the gaseous material streaming out from the B0 component, where the H and He I absorption lines originate.

(b) A region where the broad emission at 4560 Å originates.

(c) A region of diluted radiation that moves with a velocity of nearly -180 km s^{-1} .

(d) A layer where the expansion velocities reach values of the order of 1000 km s^{-1} and the resonance lines of Si IV and N IV form.

(e) A triple nebulosity with an electron density of the order of 10^4 cm⁻³ or perhaps even larger, and velocities of -18, +9 and +48 km s⁻¹. The observational evidence seems to support the idea that we are dealing with a sort of protoplanetary nebula that may be an (early?) episode in the rapid mass loss stage, and the question is open as to whether the RY Sct system is in its second episode of mass loss.

In regard to the lines that indicate dilution effect, we may say that, according to Struve & Wurm's (1938) figures, the diluted layer where the strong and narrow absorption of He I 3889 originates is probably located at a distance of more than 5 stellar radii from the star's centre. This layer is moving away from the system at a velocity of about -180 km s^{-1} , and its density is probably not higher than 10^{11} cm^{-3} or so.

As for the triple nebulosity suggested by the narrow, strong E1 emissions, we are certainly dealing with a low excitation, low density nebula because of the presence of He I in the spectrum, because of the derived value of $N_{\rm e}$ (cf. 4.1.1), and because of the steepness of the Balmer decrement (cf. Miyamoto 1952). The three nebular layers move with velocities of +48, +9 and -18 km s⁻¹, respectively (§ 4.1.1). In planetary nebulae it is common to find multiple shell structures (e.g., Chu 1989). However, the question regarding the order in which the three nebulosities are located relative to the progenitor cannot be answered at this stage.

It is possible that we are dealing with a protoplanetary nebula. The fact that the characteristics of the infrared excess of the object suggest emission by silicate grains (cf. Aitken et al. 1979; Scott 1983) and the fact that RY Sct is an *IRAS* source further support this possibility. An additional argument in favor of this hypothesis is provided by Kwok's (1990) "infrared sequence in the late stages of stellar evolution".

On the other hand, the broader H and He I absorptions must come from a region where the gas stream is important. The larger width, relative to Si IV and similarly behaved lines, arising from turbulent motions, characterizes this region. The larger eccentricity and smaller amplitude of the velocity curve may be the result of the stream and the gaseous envelope that surrounds the star.

The results suggest that the star that is presently losing mass is the less massive component of the system, the companion displaying a thick disk around it. Therefore, RY Sct is a system in the rapid massloss stage after the mass-ratio reversal has taken place. The question then arises as to whether the sort of protoplanetary nebula displayed would perhaps only be an early episode immediately after the rapid mass-loss stage.

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